
Hamiltonizing Nonholonomic Mechanics

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Nonholonomic Systems



$$L, \quad a_{ij}\dot{q}^j = 0.$$

Lagrange
d'Alembert

$$\int_a^b (\delta L + \lambda_i a_{ij} \delta q^j) dt = 0$$

Physical
Equations
of Motion

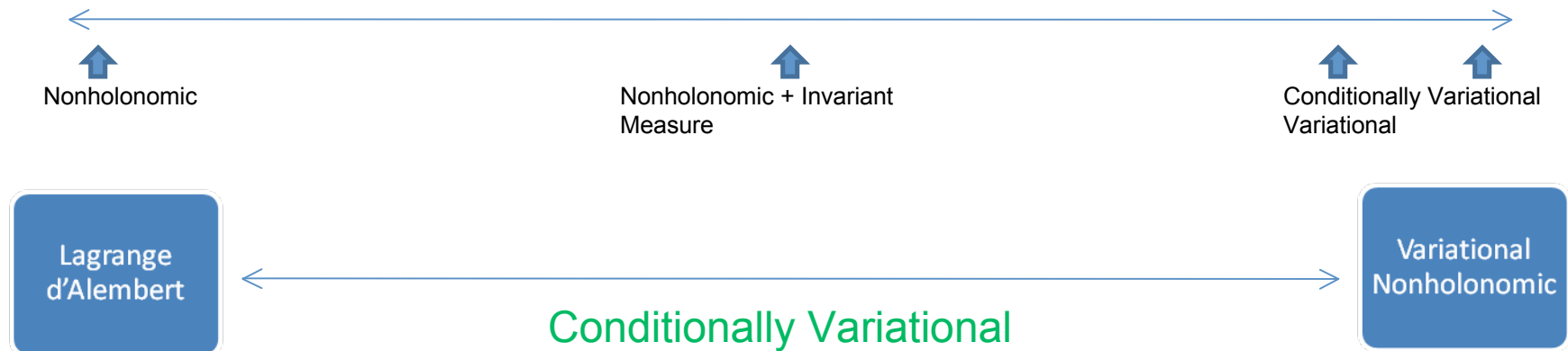
Variational
Nonholonomic

$$\delta \int_a^b (L + \lambda_i a_{ij} \dot{q}^j) dt = 0$$

Related to
Optimal
Control



A link between the two



Under the equivalence conditions and certain regularity conditions the Euler-Lagrange equations of the variational Lagrangian L_V defined by

$$L_V = L - \frac{\partial L}{\partial \dot{s}^a} (\dot{s}^a + A_\alpha^a(r) \dot{r}^\alpha)$$

reproduce the nonholonomic mechanics when the initial velocities \dot{q}_0 lie in \mathcal{D} .

Moreover, for systems with an invariant measure, such conditions depend explicitly on the invariant measure density N .



Example: Vertical Rolling Disk

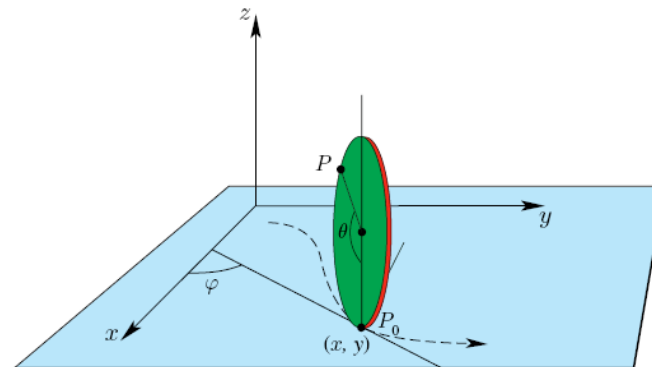
$$L = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) + \frac{1}{2}I\dot{\theta}^2 + \frac{1}{2}J\dot{\varphi}^2, \quad \dot{x} = R \cos \varphi \dot{\theta}, \quad \dot{y} = R \sin \varphi \dot{\theta}$$

Lagrange
d'Alembert



Variational
Nonholonomi
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$$\begin{aligned} \ddot{\varphi} &= 0, \\ \ddot{\theta} &= 0, \\ \dot{x} &= R \cos(\varphi) \dot{\varphi}, \\ \dot{y} &= R \sin(\varphi) \dot{\varphi}. \end{aligned}$$



$$\begin{aligned} \ddot{\varphi} &= -\frac{mR}{J} [c_1 \sin(\varphi) - c_2 \cos(\varphi)] \dot{\theta}, \\ \ddot{\theta} &= \frac{mR}{(I + mR^2)} [c_1 \sin(\varphi) - c_2 \cos(\varphi)] \dot{\varphi}, \\ \dot{x} &= R \cos(\varphi) \dot{\varphi} + c_1, \\ \dot{y} &= R \sin(\varphi) \dot{\varphi} + c_2. \end{aligned}$$

$$L_V = -\frac{1}{2}m(\dot{x}^2 + \dot{y}^2) + \frac{1}{2}I\dot{\theta}^2 + \frac{1}{2}J\dot{\varphi}^2 + mR\dot{\theta}(\dot{x}\cos \varphi + \dot{y}\sin \varphi)$$