The Existence of Maximal Slicings in Asymptotically Flat Spacetimes

M. Cantor*, A. Fischer** ***, J. Marsden⁺ ***, N. Ō Murchadha⁺⁺, and J. York⁺⁺

Department of Physics and Astronomy, University of North Carolina, Chapel Hill, North Carolina 27514, USA

Abstract. We consider Cauchy data (g, π) on \mathbb{R}^3 that are asymptotically Euclidean and that satisfy the vacuum constraint equations of general relativity. Only those (g, π) are treated that can be joined by a curve of "sufficiently bounded" initial data to the trivial data $(\delta, 0)$. It is shown that in the Cauchy developments of such data, the maximal slicing condition $\operatorname{tr} \pi = 0$ can always be satisfied. The proof uses the recently introduced "weighted Sobolev spaces" of Nirenberg, Walker, and Cantor.

Consider the set \mathscr{C} of spacetimes which are the Cauchy developments of initial data (g, π) on \mathbb{R}^3 which are asymptotically Euclidean and which satisfy the constraint equations [see (3) and (4) below] in the dynamical formulation of general relativity [1]. In 1968, Brill and Deser [2] conjectured that one can maximally slice any such spacetime, i.e. one can find spacelike hypersurfaces on which tr $\pi=0$. In a Hamiltonian analysis of general relativity tr π assumes the role of a gauge variable (see for example [12]) and so one would expect that the tr $\pi=0$ condition can be met in any such spacetime. Here we prove that the Brill-Deser conjecture is true.

We consider only those (g, π) which can be joined by a curve of "sufficiently bounded" initial data (to be explained later) to flat space $(\delta, 0)$. Thus we are considering the component \mathscr{C}_0 of $(\delta, 0)$ in the set of asymptotically Euclidean solutions of the constraint equations. \mathscr{C}_0 is restricted to those 3-metrics which are derived from Lorentz metrics on \mathbb{R}^4 that are near the "background" Minowski metric. The set \mathscr{C}_0 is discussed in [7–11].

^{*} Present address: Department of Mathematics, Duke University, Durham, North Carolina 27706, USA

^{**} Present address: Department of Mathematics, University of California, Santa Cruz, California 95064, USA

 ^{***} Research partially supported by National Science Foundation Grants GP-39060 and GP-15735
 * Present address: Department of Mathematics, University of California, Berkeley, California 94720,

USA

⁺⁺ Research partially supported by National Science Foundation Grant GP-43909 to the University of North Carolina

In her note [6], Choquet-Bruhat proves a theorem for spacetimes with compace spacelike hypersurfaces which is similar to our step 2 below. She also notes her proof extends to yield the local result for spacetimes with noncompact spacelike hypersurfaces. The authors became aware of [6] after the present work was completed.

M. Cantor et al

We shall prove:

Theorem. Let $(g, \pi) \in \mathscr{C}_0$. Then in the Cauchy development of (g, π) there is a slice on which the trace of the second fundamental form is zero. (Recall that this entails $tr\pi = 0$.

There is a similar theorem for the component of & containing any given (g,π) with $tr\pi=0$ or in the case of compact hypersurfaces, $tr\pi/\mu_g = constant$ (see [6] and [12]). The constant depends on the hypersurface. This theorem is proven similarly to the one in this paper.

The proof requires the use of the weighted Sobolev spaces $M_{s,s}^p$ introduced in [3]. For compact hypersurfaces, the usual Sobolev spaces $W^{s,p}$ will do, as in [9].

Definition. Let $\sigma(x) = (1 + |x|^2)^{1/2}$. For $1 \le p \le \infty$, s a nonnegative integer, and $\delta \in \mathbb{R}$, let $M_{1,0}^{\rho}(\mathbb{R}^{n}, \mathbb{R}^{q})$ be the completion of $C_{0}^{\infty}(\mathbb{R}^{n}, \mathbb{R}^{q})$ with respect to the norm

$$\left\|f\right\|_{p,a,\delta} = \sum_{|a| \leq s} \left(\left\|(D^{a}f)\sigma^{\delta+|a|}\right\|_{L^{p}}\right).$$

The elementary properties of these spaces are discussed in [3,4]. The important technical result for this paper is

Lemma 1. [5]. Let n > m and $A_{\infty} = \sum_{\substack{|x| = m \\ n \neq m}} \overline{a_x} D^x$ be an elliptic homogeneous operator on \mathbb{R}^n . Suppose we have an elliptic operator $A = \sum_{\substack{|x| \le m \\ n \neq m}} a_x(x) D^x$ on \mathbb{R}^n satisfying for $s \ge m$, $a_a \in \mathbb{C}^{-m}$ and

 $\sup |D'(a_{\mathbf{a}}(\mathbf{x})) \cdot \sigma^{\mathbf{m}-\mathbf{e}+|\mathbf{y}|} < \infty \quad for \quad |\mathbf{x}| < m$

and

 $\limsup |D^{\gamma}(a_{s}(x) - \bar{a}_{s})\sigma^{1\gamma i}| < \varepsilon \quad for \quad |\alpha| = m$

and $|\gamma| \leq s - m$. Then if p > n/(n-m) und $0 \leq \delta < -m + n(p-1)/p$, and c is sufficiently small, A is an isomorphism between $M_{h,h}^{p}$ and $M_{h,m,h,m}^{p}$.

Remark. The smoothness condition of the a, may be relaxed by taking completions in the appropriate Banach space of linear operators. This fact is used implicitly below.

We shall apply Lemma 1 where A is the Laplacian with respect to some asymptotically flat metric on \mathbb{R}^3 . Thus n=3 and m=2. We assume p and δ are as in the theorem and s > n/p + 2. The (g, π) we shall consider will be of the form $g = \delta + h$ with $h \in M_{s,\delta}^p$ and $\pi \in M_{s-1,\delta+1}^p$ (see [8]). All norms are taken with respect to the flat background metric. Note that for $g \in \mathscr{C}_0$ these norms are equivalent to those induced by g. Note we may take $tr \pi \in M_{i-1,i+2}^{p}$. The topology on the space of initial data is given by the $M_{4,4}^p$ norms.

g : The required slicing w Letting the shift vector is a density): in the second $\partial g/\partial l = 2N(\pi - \frac{1}{2})(tr\pi)$ $\partial \pi/\partial \lambda = -N(\operatorname{Ric}(q))$ $-2N(\pi \times \pi$ $\mathcal{L}(g,\pi) = (\pi \cdot \pi - \frac{1}{2})(tr$ $\delta_n \pi = 0$

"Maximal Slicings."

and using $p = tr \pi / \mu_{e} = 2$

 $-\partial p/\partial \lambda = 2(k \cdot k - P^2)N$

Step. 1. If p=0 for some which the dynamics is do

Proof. Writing $N = 1 + \hat{N}$

 $\Rightarrow \partial p/\partial \lambda = 2k \cdot k + 2(k \cdot k -$

Thus the equation $\partial p/\partial c$. for each 1. Thus for this will be maintained.

In what follows we : p=0 at $\lambda=1$ by choosing

Step 2. (Local Argument neighborhood V of $(g_0,$ such that p=0 at $\lambda=1$. (F reached by the dynamics.

Proof. Let $F = \mathscr{C}_0 \times \mathscr{M}_{**}$ open set) by -*

 $F((g, \pi), N) = \{$ the func

Then using smoothne a smooth mapping. The d δN is

 $D_N F((g_0, \pi_0), 0) \cdot \delta N = |$

where $k_0(\lambda)$ is the evolution for $g_0(\lambda)$.

Since we are only coneasily from Lemma 1 th Thus by the implicit fur for $N(g, \pi)$ near 0 and (g, π)

M. Cantor et al.

or spacetimes with compact 2 below. She also notes her with noncompact spacelike fter the present work was

tent of (g, π) there is a slice tro. (Recall that this entails

of \mathscr{C} containing any given surfaces, tr $\pi/\mu_g =$ constant persurface. This theorem is

lev spaces $M_{1,\delta}^{\mu}$ introduced ev spaces $W^{s,\mu}$ will do, as

1 nonnegative integer, and 1) with respect to the norm

cussed in [3, 4].

ic homogeneous operator
a_(x)D^{*} on R^{*} satisfying

(p-1)/p, and ε is sufficiently

laxed by taking completions This fact is used implicitly

scian with respect to some =2. We assume p and δ are consider will be of the form forms are taken with respect these norms are equivalent 1.4+2. The topology on the Maximal Slicings

The required slicing will be determined by a lapse function $N(\lambda, x) = (-g^{00})^{-1/2}$. Letting the shift vector $g_{0i} = X_i(\lambda, x) = 0$, the Einstein Equations read (here π is a density):

 $\hat{c}g/\hat{c}\lambda = 2N(\pi - \frac{1}{2}(\operatorname{tr}\pi)g)(1/\mu_{g}) = -2Nk$ (1) $\hat{c}\pi/\hat{c}\lambda = -N(\operatorname{Ric}(g) - \frac{1}{2}R(g)g)\mu_{g} + \frac{1}{2}N(\pi \cdot \pi - \frac{1}{2}(\operatorname{tr}\pi)^{2})/\mu_{g} -2N(\pi \times \pi - \frac{1}{2}(\operatorname{tr}\pi)\pi)/\mu_{g} + (\operatorname{Hess} N - gV^{2}N)\mu_{g}$ (2)

 $\mathcal{H}(g,\pi) = (\pi \cdot \pi - \frac{1}{2} (\mathrm{tr} \pi)^2) / \mu_g - R(g) \mu_g = 0$ $\delta_o \pi = 0$ (3)

v_g, *v* = *v* (4)

and using $p = \operatorname{tr} \pi/\mu_{\phi} = 2 \operatorname{tr} k$, we find from the above equations that

$$\tilde{c}p/\tilde{c}\lambda = 2(k\cdot k - V^2)N.$$
⁽⁵⁾

Step. 1. If p=0 for some λ , we may choose an N such that p is zero for all λ (for which the dynamics is defined).

Proof. Writing $N = 1 + \tilde{N}$ (so that \tilde{N} is close to 0 when N is close to 1), we find $\tilde{c}p/\tilde{c}\lambda = 2k \cdot k + 2(k \cdot k - F^2)\tilde{N}$.

Thus the equation $\partial p/\partial \lambda = 0$ may be solved using Lemma 1 for $\tilde{N}(\lambda) \in M_{s+1,s}^p$ for each λ . Thus for this choice of $N = 1 + \tilde{N}$ in the dynamics the condition p = 0 will be maintained. \Box

In what follows we show that whatever p equals at $\lambda = 0$, we may achieve p=0 at $\lambda = 1$ by choosing a suitable N. Throughout, we shall take $\partial N/\partial \lambda = 0$.

Step 2. (Local Argument). Let $(g_0, \pi_0) \in \mathscr{C}_0$ and suppose $\operatorname{tr} \pi_0 = 0$. Then there is a neighborhood V of (g_0, π_0) such that if $(g, \pi) \in V$ then there is an $N \in \mathscr{M}_{s+1,\delta}^p$ such that p=0 at $\lambda=1$. (By a suitable choice of scale, we may assume $\lambda=1$ will be reached by the dynamics.)

Proof. Let $F = \mathscr{C}_0 \times \mathscr{M}^p_{s+1,\delta}(\mathbb{R}^3, \mathbb{R}) \to \mathcal{M}^p_{s-1,\delta+2}(\mathbb{R}^3, \mathbb{R})$ be defined (on a suitable open set) by

 $F((g, \pi), N) = \{$ the function p at $\lambda = 1$ determined by Equations (1), (2), (5) $\}$.

Then using smoothness properties of the evolution equations (see [8]), F is a smooth mapping. The derivative with respect to N at $((g_0, \pi_0), 0)$ in the direction δN is

$$D_N F((g_0, \pi_0), 0) \cdot \delta N = \left(\int_0^1 (k_0(\lambda) \cdot k_0(\lambda) - \mathcal{V}_{\lambda}^2) d\lambda\right) \delta N$$

(6)

where $k_0(\lambda)$ is the evolution of k_0 for the given (g_0, π_0) and V_{λ}^2 is the Laplacian for $g_0(\lambda)$.

Since we are only considering functions that are independent of λ , it follows easily from Lemma 1 that the operator (6) is an isomorphism (see also [9]). Thus by the implicit function theorem we can uniquely solve $F((g, \pi), N) = 0$ for $N(g, \pi)$ near 0 and (g, π) near (g_0, π_0) . This proves step 2.

.....

189

Step 3 (Globalization). Let (g_0, π_0) be joined to (g, π) be a continuous curve $(g(\alpha), \pi(\alpha))$ in $\mathscr{C}_0, \alpha \in [0, 1]$. Let J be the set of α for which the resulting space time has a maximal slice. Then $0 \in J$ and step 2 shows that J is open. We can always work in a neighborhood of the curve $(g(\alpha), \pi(\alpha))$ so that the evolution times used in step 2, can be chosen to be uniform along the curve.

To show J is closed, let $\alpha_m \in J$ and $\alpha_m \to \alpha$. Let N_m be the unique lapse functions given by step 2. In order to demonstrate that \mathcal{P}_m^2 remains uniformly elliptic and the slices "uniformly spacelike", we may take a sequence of coordinate transformations f_m on the slices S_m chosen so as to keep the eigenvalues of g_m (relative to the flat background metric) bounded away from zero. Since $k_0(m)$ remains uniformly bounded and \mathcal{P}_m^2 remains uniformly elliptic for $\alpha_m, m \to \infty$, the N_m will converge to a function N. This N is the required zero of F.

Thus J = [0, 1] and our proof is complete.

Note Added in Prixif. The hypotheses of Lemma 1 should include that a_0 is non-positive. In our application, $a_0 = -\int_{1}^{1} k_0(\lambda) \cdot k_0(\lambda) d\lambda \leq 0$.

References

190

- 1. Arnowitt, R., Deser, S., Misner, C. W.: The dynamics of general relativity, In: Gravitation: an introduction to current research, (ed. L. Witten). New York: Wiley 1962
- Brill, D., Deser, S.: Variational methods and positive energy in general relativity. Ann. Phys. 50, 548-570 (1968)
- 3. Cantor, M.: Spaces of functions with asymptotic conditions on R⁴. Ind. U. J. Math. to appear (1975)
- 4. Cantor, M.: Perfect fluid flows over R* with asymptotic conditions, J. Funct. Anal. 18. 73-84 (1975)
- Cantor, M.: Growth of Solutions of elliptic equations with nonconstant coefficients on R^a. Preprint
 Choquet-Bruhat, Y.: Sous-Varietes maximales, ou a courbure constante, de varietes lorentziennes, C. R. Acad. Sc. Paris, 280, Ser A, 169-171 (1975)
- Fischer, A., Marsden, J.: The Einstein equations of evolution A geometric approach. J. Math. Phys. 13, 546-568 (1972)
- Fischer, A., Marsden, J.: The Einstein evolution equations as a first order quasi-linear hyperbolic system. Commun. math. Phys. 28, 1-38 (1972)
- Fischer, A., Marsden, J.: Linearization stability of nonlinear partial differential equations, Proc. Symp. Pure Math. A. M. S. 27, 219-263 (1975) (also Bull. A. M. S. 79, 997-1003 (1973) 80, 479-484, and General Relativity and Gravitation 5, 73-77 (1974)
- 10. O'Murchadha, N., York, J. W.: Initial value problem of general relativity (1, 11). Phys. Rev. D 10, 428-436, 437 -446 (1974)
- 11. O'Murchadha, N., York, J.W.: Gravitational energy. Phys. Rev. D 10, 2345-2357 (1974)
- 12. York, J.W.: The role of conformal three geometry in the dynamics of gravitation. Phys. Rev. Letters 28, 1082 (1972)

Communicated by J. Ehlers

Received July 9, 1975

Manuscripts should be written in English; paper Manuscripts must be in t with double spacing an whenever possible. Mirr. legible.

INS

M. Cantor et al.

Please include a "Note suggestion overleaf.

To speed up publication sionally numbered page graphical errors only; the additions or deletions to Diagrams should be sub-They should be drawn in twice the size of the fini figure 1, for example, to for reduction $\times \frac{1}{2}$). The a where diagrams may be i Footnotes, other than the bered consecutively and pi the end of the article). Please give on the first par which should not exceed References to the literatu The following information initials of all authors, n numbers and year of F name(s) of author(s), full year of publication.

Bombieri, E., Giusti, E.: In Tate, J.T.: p-Divisible grfields, pp. 158-183, Ber: